

POLAR COORDINATES FOR THE GENERALIZED BAOUENDI-GRUSHIN OPERATOR AND APPLICATIONS*

Dou Jingbo, Niu Pengcheng and Han Junqiang

(Department of Applied Mathematics, Northwestern Polytechnical University,
Xi'an 710072, China)

(E-mail: djbmn@126.com(dou), pengchengniu@yahoo.com.cn(Niu))

(Received May 26, 2006)

Abstract In this paper, by using the polar coordinates for the generalized Baouendi-Grushin operator

$$\mathcal{L}_\alpha = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2} + \sum_{j=1}^m |x|^{2\alpha} \frac{\partial^2}{\partial y_j^2},$$

where $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n, y = (y_1, y_2, \dots, y_m) \in \mathbb{R}^m, \alpha > 0$, we obtain the volume of the ball associated to \mathcal{L}_α and prove the nonexistence for a second order evolution inequality which is relative to \mathcal{L}_α .

Key Words Generalized Baouendi-Grushin operator; polar coordinate; nonexistence; second order evolution inequality.

2000 MR Subject Classification 35R45, 35J60.

Chinese Library Classification O175.2.

1. Introduction

The polar coordinates for the Heisenberg group \mathbb{H}_1 and for the Heisenberg group \mathbb{H}_n were defined by Greiner [1] and D'Ambrosio [2], respectively. In [3] and [4], such coordinates for the Grushin operator in \mathbb{R}^{n+1} and generalized Baouendi-Grushin operator

$$\mathcal{L}_\alpha = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2} + \sum_{j=1}^m |x|^{2\alpha} \frac{\partial^2}{\partial y_j^2} \tag{1.1}$$

were studied, where $x = (x_1, x_2, \dots, x_n) \in \mathbb{R}^n, y = (y_1, y_2, \dots, y_m) \in \mathbb{R}^m, \alpha > 0$.

Nonexistence results of positive solutions for singular elliptic inequality, parabolic and hyperbolic inequality in the Euclidean space \mathbb{R}^n have been largely considered, see

*The project supported by Natural Science Basic Research Plan in Shanxi Province of China, Program No.2006A09

[5, 6] and their references. The singular sub-Laplace inequality and related evolution inequalities in the Heisenberg group \mathbb{H}_n were studied in [2, 7]. The singular sub-elliptic inequality and the first order evolution inequality related to \mathcal{L}_α were considered in [4].

In this paper we will give some applications of the polar coordinates for the operator \mathcal{L}_α . In particular, we explicitly compute the volume of the ball in the sense of the distance associated with \mathcal{L}_α . Also, we discuss the nonexistence for a second order evolution inequality which is relative to \mathcal{L}_α

$$\begin{cases} u_{tt} - \frac{d^2}{\psi_{2\alpha}} \mathcal{L}_\alpha(au) \geq |u|^q, & \text{on } \mathbb{R}_*^{n+m} \times (0, +\infty), \\ u(x, y, 0) = u_0(x, y), & \text{on } \mathbb{R}_*^{n+m}, \\ u_t(x, y, 0) = u_1(x, y), & \text{on } \mathbb{R}_*^{n+m}, \end{cases} \quad (1.2)$$

where $a \in L^\infty(\mathbb{R}^{n+m} \times [0, +\infty))$, $q \geq 1$, $\mathbb{R}_*^{n+m} = \mathbb{R}^{n+m} \setminus \{(0, 0)\}$. Our consideration is motivated by D'Ambrosio [2]. Let us note that some essential differences appearing unavoidably. We recall some known facts about the operator \mathcal{L}_α (see[8]). Let

$$Z_i = \frac{\partial}{\partial x_i}, \quad Z_{n+j} = |x|^\alpha \frac{\partial}{\partial y_j} \quad (i = 1, 2, \dots, n; j = 1, 2, \dots, m). \quad (1.3)$$

Denote the generalized gradient

$$\nabla_L = (Z_1, \dots, Z_n, Z_{n+1}, \dots, Z_{n+m}).$$

There exists a natural family of anisotropic dilations attached to \mathcal{L}_α , i.e.,

$$\delta_\lambda(x, y) = (\lambda x, \lambda^{\alpha+1} y), \quad \lambda > 0, (x, y) \in \mathbb{R}^{n+m}.$$

It leads to a homogeneous dimension for \mathcal{L}_α

$$Q = n + (\alpha + 1)m.$$

One easily examines that

$$\mathcal{L}_\alpha \circ \delta_\lambda = \lambda^2 \delta_\lambda \circ \mathcal{L}_\alpha,$$

so that \mathcal{L}_α becomes homogeneous of degree two with respect to the anisotropic dilations.

Introduce the distance function

$$d(x, y) = (|x|^{2(\alpha+1)} + (\alpha + 1)^2 |y|^2)^{\frac{1}{2(\alpha+1)}}. \quad (1.4)$$

It should also be noted that

$$|\nabla_L d|^2 = \psi_{2\alpha} = \frac{|x|^{2\alpha}}{d^{2\alpha}} \quad (1.5)$$

and

$$\mathcal{L}_\alpha d = \psi_{2\alpha} \frac{Q-1}{d}. \quad (1.6)$$