
THE CAUCHY PROBLEM OF NONLINEAR SCHRÖDINGER-BOUSSINESQ EQUATIONS IN $H^s(R^d)$

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Abstract In this paper, the local well posedness and global well posedness of solutions for the initial value problem (IVP) of nonlinear Schrödinger-Boussinesq equations is considered in $H^s(R^d)$ by resorting Besov spaces, where real number $s \geq 0$.

Key Words Schrödinger-Boussinesq equation; global solutions in Besov spaces.

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1. Introduction

We consider the existence and uniqueness of the local solutions and global solutions for the following initial value problem (IVP) of nonlinear Schrödinger-Boussinesq equations

$$i\epsilon_t + \Delta\epsilon - n\epsilon - A|\epsilon|^p\epsilon = 0, \quad (1.1)$$

$$n_{tt} - \Delta(n - \Delta n + Bn^{K+1} + |\epsilon|^2) = 0, \quad x \in R^d, t \in R, \quad (1.2)$$

$$\epsilon(x, 0) = \epsilon_0(x), \quad n(x, 0) = n_0(x), \quad n_t(x, 0) = \Delta\phi_0(x), \quad x \in R^d, \quad (1.3)$$

where A and B are constants, K is a positive integer, real number $p > 0$; ϵ and ϵ_0 are complex functions; n , n_0 and ϕ_0 are real functions; Δ is Laplacian operator in R^d .

The nonlinear Schrödinger (NLS) equation models a wide range of physical phenomena including self-focusing of optical beams in nonlinear media, propagation of Langmuir waves in plasmas, etc. (see [1] and the references therein). Boussinesq equation as a model of long waves is derived in the studies of the propagation of long waves on the surface of shallow water[2], the nonlinear string [3] and the shape-memory alloys[4], etc. The nonlinear Schrödinger-Boussinesq equations (1.1)(1.2) is considered as a model of interactions between short and intermediate long waves, which is derived

in describing the dynamics of Langmuir soliton formation and interaction in a plasma [5-7] and diatomic lattice system [8], etc.

The Solitary wave solutions and integrability of nonlinear Schrödinger-Boussinesq equations has been considered by several authors, see [5, 6, 9] and the references therein. In [10] Guo established the existence and uniqueness of global solution for IVP (1.1)–(1.3) in H^k (integer $k \geq 4$) with $d = 1$ and $A = 0$. In [11] the existence and uniqueness of global solution for Cauchy problem of dissipative Schrödinger-Boussinesq equations in H^k (integer $k \geq 4$) with $d = 3$ is proved by Guo and Shen. For damped and dissipative Schrödinger-Boussinesq equations with initial boundary value, the existence of global attractors and the finiteness of the Hausdorff and the fractal dimensions of the attractor is established by Guo and Chen ([12], $d=1$) and Li and Chen ([13], $d \leq 3$), respectively.

In this paper, the local well-posedness in H^s , the conservation of energy and the global well-posedness in H^s (real number $s \geq 1$ and $d = 1, 2, 3$) of IVP (1.1)–(1.3) is proved.

Definition 1(admissible pair) *The pair (q, r) is admissible if $\frac{2}{q} = d(\frac{1}{2} - \frac{1}{r})$;*

$2 \leq r \leq \infty$ for $d = 1$, $2 \leq r \leq \infty$ for $d = 2$, $2 \leq r < \frac{2d}{d-2}$ for $d \geq 3$.

Definition 2(condition $P(m)$) *For a positive integer m , it is called that p satisfies the condition $P(m)$ if either p is an even integer, or p is not an even integer and $p + 1 > m$.*

The main theorems of this paper are stated as follows.

Theorem 1 *Suppose that $\epsilon_0, n_0, \phi_0 \in H^s(R^d)$, $0 \leq s < \frac{d}{2}$, K is an integer, p satisfies the condition $P([s] + 1)$, $0 < p$, $K \leq \frac{4}{d-2s}$; then for any admissible pair (q, r) , there exists $T = T(\epsilon_0, n_0, \phi_0) > 0$ and a unique solution (ϵ, n) of IVP (1.1)–(1.3) such that*

$$\epsilon, n, (-\Delta)^{-1}n_t \in L^q(0, T; B_{r,2}^s(R^d)) \cap C([0, T]; H^s(R^d))$$

Moreover, this solution has the following additional properties.

(I) *Let $p, K < \frac{4}{d-2s}$. If $\epsilon_{0j}, n_{0j}, \phi_{0j}$ are sequences in $H^s(R^d)$ with $(\epsilon_{0j}, n_{0j}, \phi_{0j}) \rightarrow (\epsilon_0, n_0, \phi_0)$, then there exists $\tilde{T} = \tilde{T}(\epsilon_0, n_0, \phi_0) \in (0, T]$, such that the solutions $(\epsilon_j, n_j) \rightarrow (\epsilon, n)$ and $(-\Delta)^{-1}\partial_t n_j \rightarrow (-\Delta)^{-1}n_t$ in $L^q(0, \tilde{T}; L^r(R^d))$, where (ϵ_j, n_j) are solutions of IVP (1.1)–(1.3) with $(\epsilon_0, n_0, \phi_0)$ replaced by $(\epsilon_{0j}, n_{0j}, \phi_{0j})$. If $s \geq 1$, then $(\epsilon_j, n_j) \rightarrow (\epsilon, n)$ and $(-\Delta)^{-1}\partial_t n_j \rightarrow (-\Delta)^{-1}n_t$ in $C([0, \tilde{T}]; H^{s-1}(R^d)) \cap L^q(0, \tilde{T}; B_{r,2}^{s-1})$. Moreover, if p satisfies the condition $P([s] + 2)$, then $(\epsilon_j, n_j) \rightarrow (\epsilon, n)$ and $(-\Delta)^{-1}\partial_t n_j \rightarrow (-\Delta)^{-1}n_t$ in $C([0, \tilde{T}]; H^s(R^d)) \cap L^q(0, \tilde{T}; B_{r,2}^s)$.*

(II) *There exists $T^* = T^*(\epsilon_0, n_0, \phi_0) > 0$ such that the solution $\epsilon, n, (-\Delta)^{-1}n_t \in C([0, T^*]; H^s(R^d)) \cap L_{loc}^q(0, T^*; B_{r,2}^s(R^d))$. If $T^* < \infty$, then*

$$\lim_{t \rightarrow T^*} \left\{ \|(-\Delta)^{\frac{s}{2}} \epsilon(\cdot, t)\|_{L^2} + \|(-\Delta)^{\frac{s}{2}} n(\cdot, t)\|_{L^2} + \|(-\Delta)^{\frac{s-2}{2}} n_t(\cdot, t)\|_{L^2} \right\} = +\infty.$$