## The Negative Spectrum of Schrödinger Operators with Fractal Potentials

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**Abstract.** Let  $\Gamma \subset \mathbb{R}^2$  be a regular anisotropic fractal. We discuss the problem of the negative spectrum for the Schrödinger operators associated with the formal expression

$$H_{\beta} = id - \Delta + \beta tr_{h}^{\Gamma}, \quad \beta \in \mathbb{R},$$

acting in the anisotropic Sobolev space  $W_2^{1,\alpha}(\mathbb{R}^2)$ , where  $\Delta$  is the Dirichlet Laplanian in  $\mathbb{R}^2$  and  $tr_h^{\Gamma}$  is a fractal potential (distribution) supported by  $\Gamma$ .

**Key Words**: Anisotropic function space, anisotropic fractal, Schrödinger operators, negative eigenvalues.

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## 1 Introduction

Let a(x,D) be a positive-definite, self-adjoint pseudodifferential operator in  $L_2(\mathbb{R}^2)$ ,  $W_1^2(\mathbb{R}^2)$  the classical Sobolev space. In general, we studied the operators of the type

$$H_{\beta} = a(x, D) + \beta V, \quad \beta \in \mathbb{R}, \tag{1.1}$$

where *V* is a generalized potential given by a distribution supported by a set  $\Gamma \subset \mathbb{R}^2$  with Lebesgue measure  $|\Gamma| = 0$ . Of interest is the number of eigenvalues, counted with respect to their multiplicities, in  $(-\infty, 0]$  in dependence on  $\beta \to \infty$ . Recall that  $\sharp M$  denotes the number of a finite set *M*. Let  $\sigma(H_\beta)$  be the spectrum of self-adjoint operator  $H_\beta$  in  $L_2(\mathbb{R}^2)$ , then

$$\sharp \{ \sigma(H_{\beta}) \cap (-\infty, 0] \} \leq \sharp \{ k \in \mathbb{N} : \sqrt{2}\beta e_k \left( a(x, D)^{-1}V \right) \geq 1 \}.$$

$$(1.2)$$

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The study of  $\sharp \{ \sigma(H_{\beta}) \cap (-\infty, 0] \}$  is sometimes called a problem of the negative spectrum.

The interest of the problem comes from quantum mechanics with the typical hydrogenlike operator

$$H = -\frac{h^2}{2m}\Delta + V(x), \qquad (1.3)$$

where *h* is the Planck's constant, *m* is mass of the electron and V(x) is the potential. Asking for (negative) eigenvalues one can divide *H* by  $h^2/2m$  and obtains

$$H_{\beta} = -\Delta + \beta V(x), \tag{1.4}$$

where  $\beta$  is proportional to  $h^2$ . Of interest is the semi-classical limit where  $h \rightarrow 0$ , or,  $\beta \rightarrow \infty$ , hence the behaviour of the cardinal number in dependence on  $\beta$ .

In slight modification of (1.3), we are interested in the negative spectrum of

$$H_{\beta} = id - \Delta + \beta t r_b^{\Gamma}, \qquad (1.5)$$

where  $tr_b^{\Gamma}$  is the interpretation of  $btr_{\Gamma}$  and  $b \in L_r(\Gamma)$  is real with  $0 \le 1/r < 1$ ,  $tr_{\Gamma}$  is the trace operator. Let  $N_{\beta}$  be the number of negative eigenvalues. If  $\Gamma$  is an (isotropic) *d*-set, we obtained satisfactory results,  $N_{\beta} \sim \beta^1$ .

For anisotropic fractals in  $\mathbb{R}^2$ , an anisotropic *d*-set  $\Gamma$  having deviation,  $0 \le a \le 1$  is, roughly speaking, a compact set which can be covered for any  $j \in \mathbb{N}_0$  with  $N_j \sim 2^{jd}$  rectangles  $\{R_l^j : l = 1, \dots, N_j\}$ ,  $c2^{-2j} \le volR_l^j \le 2^{-2j}$ , for some c, 0 < c < 1, independent of j and l, having the sides parallel to the axes of coordinates and side lengths  $2^{-a_1^{j,l}}$ ,  $2^{-a_2^{j,l}}$  satisfying

$$(1-a)j \le a_1^{j,l} \le a_2^{j,l} \le (1+a)j$$

for any  $l = 1, \dots, N_i$ . In this case, one has for  $N_\beta$  only two side estimates of type

$$C_1 \beta^{\omega_1} \le N_\beta \le C_2 \beta^{\omega_2} \tag{1.6}$$

for appropriate positive numbers  $\omega_1$  and  $\omega_2$  with  $\omega_1 \leq 1 \leq \omega_2$ .

The aim of this paper is to show that sharp estimate of (1.6), restricting ourselves to the regular anisotropic fractals and the Schrödinger operators acting in the anisotropic Sobolev space  $W_2^{1,\alpha}(\mathbb{R}^2)$ . We will show that

$$N_{\beta} \sim \beta^1. \tag{1.7}$$

In Section 2 we present basic facts concerning anisotropic function spaces, regular anisotropic fractals and some preparatory facts for the proofs. The main results, containing the precise proof of (1.6), are presented in Section 3.

382