

A New Higher Order Fractional-Step Method for the Incompressible Navier-Stokes Equations

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Abstract. In this paper, we present a rigorous error analysis of a new higher order fractional-step scheme for approximation of the time-dependent Navier-Stokes equations. The main feature of the proposed scheme is twofold. First, it is a two-step scheme in which the incompressibility and nonlinearities are split. Second, this scheme is a linear scheme and is simple to implement. It is shown that the proposed scheme possesses the convergence rate $\mathcal{O}((\Delta t)^{3/2})$ in the discrete $l^2(\mathbf{H}_0^1) \cap l^\infty(\mathbf{L}^2)$ -norm for the end-of-step velocity. Two different numerical experiments are presented to confirm the theoretical analysis and the efficiency of the proposed scheme.

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Key words: Incompressible Navier-Stokes equations, fractional-step method, Crank-Nicolson scheme, temporal errors estimates.

1 Introduction

The Navier-Stokes equations are used to describe the flow of a viscous and incompressible fluid, which are governed by the following time-dependent nonlinear problems:

$$\frac{\partial \mathbf{u}}{\partial t} - \mu \Delta \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} + \nabla p = \mathbf{f}, \quad (1.1a)$$

$$\operatorname{div} \mathbf{u} = 0, \quad (1.1b)$$

for $x \in \Omega$ and $t \in (0, T]$ with $T > 0$, where Ω is an open bounded domain in \mathbf{R}^d ($d=2$ or 3) with a sufficiently smooth boundary $\partial\Omega$. The constant $\mu > 0$ represents the kinematic viscosity.

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The vector-value function \mathbf{f} represents the body forces applied to the fluid. To study (1.1a)-(1.1b), the appropriate initial and boundary conditions are needed. For the sake of simplicity, in this paper, we consider the following initial and boundary conditions:

$$\mathbf{u}(x,0) = \mathbf{u}_0 \quad \text{in } \Omega \quad \text{and} \quad \mathbf{u} = 0 \quad \text{on } \partial\Omega \times [0, T]. \quad (1.2)$$

In (1.1a)-(1.2), the unknown are the fluid velocity \mathbf{u} and the fluid kinematic pressure p which are coupled by the incompressible condition $\operatorname{div} \mathbf{u} = 0$. This condition is one of the main concern in designing efficient time discrete schemes for numerical simulation of (1.1a)-(1.2). The well-known projection method, initially proposed by Chorin [4] and Temam [22], is designed to decouple the velocity and pressure, and has been further developed in various directions [12, 19–21]. This method is first to compute an intermediate velocity field without taking into account incompressibility and then perform a pressure correction, which is a projection back to the subspace of solenoidal (divergence-free) vector field. Shen proved that this projection method is first order accurate in the time step size [19]. The incompatibility of the projection boundary conditions with (1.2) may result in a numerical boundary layer of size $\mathcal{O}(\sqrt{\Delta t})$, where Δt is the time step size [18, 24].

The viscosity-splitting fractional-step method, proposed by Blasco-Codina-Huerta [2], also is an efficient algorithm for numerical simulation of the incompressible Navier-Stokes equations. It is a two-step scheme in which the incompressibility and nonlinearities of the Navier-Stokes problems are split into different steps and allows the enforcement of the original boundary conditions in all substeps. It was shown that the intermediate and end-of-step velocities converge to a continuous solution in $L^2(\Omega)$ and $\mathbf{H}_0^1(\Omega)$ [2]. Moreover, these velocities and pressure were first-order accurate in the time step size [3, 10]. Subsequently, Dai studied a nonlinear higher order viscosity splitting, fractional-step scheme [5]. However, one has to solve a nonlinear problem at each time step, which results in a time-consuming in the practical computations. Recently, the first-order viscosity-splitting fractional-step methods have been applied to other nonlinear partial differential systems, such as the three-dimensional incompressible MHD systems [1] and the primitive equations in the field of geophysical fluids [11].

In this paper, based upon Crank-Nicolson discretization scheme in time, we will study a higher order fractional-step scheme for the approximation of (1.1a)-(1.2). Unlike the nonlinear scheme in [5], the proposed scheme is a semi-implicit scheme and it only solves two linear systems at every time-step. Therefore, it is simple to implement. To state main result, we introduce the following notations. Let X be a Banach space equipped with norm $\|\cdot\|_X$. Let $0 = t_0 < t_1 < \dots < t_N = T$ be a uniform partition of time interval $[0, T]$ with the time step size $\Delta t = T/N$ and $t_n = n\Delta t$ for $0 \leq n \leq N$. We denote two discrete norms by

$$\|\mathbf{u}^n\|_{l^2(X)} = \left(\Delta t \sum_{n=1}^N \|\mathbf{u}^n\|_X^2 \right)^{1/2}, \quad \|\mathbf{u}^n\|_{l^\infty(X)} = \max_{1 \leq n \leq N} \|\mathbf{u}^n\|_X.$$

In this paper we will show that the proposed higher order fractional-step scheme provides the temporal error estimates of $\mathcal{O}((\Delta t)^{3/2})$ for the end-of-step velocity in $l^2(\mathbf{H}_0^1) \cap$