

ON THE L_∞ CONVERGENCE AND THE EXTRAPOLATION METHOD OF A DIFFERENCE SCHEME FOR NONLOCAL PARABOLIC EQUATION WITH NATURAL BOUNDARY CONDITIONS^{*1)}

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Abstract

In paper [4] (J. Comput. Appl. Math., 76 (1996), 137-146), a difference scheme for a class of nonlocal parabolic equations with natural boundary conditions was derived by the method of reduction of order and the unique solvability and second order convergence in L_2 -norm are proved. In this paper, we prove that the scheme is second order convergent in L_∞ norm and then obtain fourth order accuracy approximation in L_∞ norm by extrapolation method. At last, one numerical example is presented.

Key words: Parabolic, Nonlocal, L_∞ convergence, Extrapolation method.

1. Introduction

Nonlocal parabolic equations have many applications. For example, in considering fluid flow in a saturated porous medium, the equation governing the pore pressure $p(r, t)$ in an annular cylindrical rock sample is given in [1] as

$$\frac{\partial p}{\partial t} + \frac{2(\nu_\mu - \nu)}{\eta(1 - \nu_\mu)} \frac{dC(t)}{dt} = c \left(\frac{\partial^2 p}{\partial r^2} + \frac{1}{r} \frac{\partial p}{\partial r} \right), \quad 0 < a < r < b, t > 0 \quad (1.1)$$

together with

$$C(t) = (b^2 - a^2)^{-1} \left[\frac{a^2}{2} p_0 + \eta \int_a^b \rho p(\rho, t) d\rho \right] \quad (1.2)$$

where c is a material constant with dimensions of velocity, the coefficient of consolidation, ν, ν_μ are shear coefficients and η a material constant, r denotes radial distance and t time. If t is replaced by ct and $\frac{2(\nu_\mu - \nu)}{\eta(1 - \nu_\mu)}$ by q in (1.1), we get

$$\frac{\partial p}{\partial t} + q \frac{dC(t)}{dt} = \frac{\partial^2 p}{\partial r^2} + \frac{1}{r} \frac{\partial p}{\partial r}, \quad 0 < a < r < b, t > 0. \quad (1.3)$$

Various initial and boundary conditions can be considered. The existence and uniqueness of the solution of (1.3) with (1.2) under some initial and boundary conditions are proved in [2] provided

$$q\eta \neq -2. \quad (1.4)$$

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It is also pointed out that, in practice, $q > 0$ and $\eta > 0$, so that (1.4) is normally met [2].

Substituting (1.2) into (1.3), we obtain an alternative form of (1.3)

$$\frac{\partial p}{\partial t} - \epsilon \int_a^b \rho \frac{\partial p}{\partial t}(\rho, t) d\rho = \frac{\partial^2 p}{\partial r^2} + \frac{1}{r} \frac{\partial p}{\partial r},$$

where

$$\epsilon = -q\eta(b^2 - a^2)^{-1}.$$

(1.4) is equivalent to $\epsilon \neq 2(b^2 - a^2)^{-1}$, or, $\frac{1}{2}(b^2 - a^2)\epsilon \neq 1$. When $q > 0$ and $\eta > 0$, we have $\epsilon < 0$. In the following, we suppose

$$\frac{1}{2}(b^2 - a^2)\epsilon < 1. \tag{1.5}$$

It is always valid when $\epsilon < 0$.

As usual, we write r by x , p by u . Lin, Tait [3] considered the finite difference solution to the nonlocal parabolic equation given by

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + \frac{1}{x} \frac{\partial u}{\partial x} + \epsilon \int_a^b \rho \frac{\partial u}{\partial t}(\rho, t) d\rho, \quad 0 < a < x < b, 0 < t \leq T, \tag{1.6}$$

subject to suitable initial and boundary conditions. If the initial-boundary conditions are of the form

$$\begin{aligned} u(x, 0) &= \phi(x), & a \leq x \leq b, \\ u(a, t) &= f(t), & u(b, t) = g(t), \quad 0 < t \leq T, \end{aligned} \tag{1.7}$$

a backward Euler scheme and a Crank-Nicolson scheme are presented, with the former giving rise to an error $O(\tau + h^2)$ and the latter to an error $O(\tau^2 + h^2)$. If the natural boundary conditions

$$\begin{aligned} u(x, 0) &= \phi(x), & a \leq x \leq b, \\ u(a, t) &= f(t), & \frac{\partial u}{\partial x}(b, t) + u(b, t) = g(t), \quad 0 < t \leq T \end{aligned} \tag{1.8}$$

are imposed, a difference scheme whose convergence order is only one in space and in time is presented. Sun [4] continually studied the finite difference solution to (1.6) with (1.8) and constructed a difference scheme by the method of reduction of order. He proved that the difference scheme is uniquely solvable and unconditionally convergent with the convergence order $O(\tau^2 + h^2)$ in energy norm. In this paper, we will prove that Sun's difference scheme is also second order convergent in L_∞ -norm and then obtain a fourth order accuracy approximation in L_∞ -norm by once extrapolation [5,6]. At last, we present a numerical example.

For generality, instead of (1.6), consider the inhomogeneous equation

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + \frac{1}{x} \frac{\partial u}{\partial x} + \epsilon \int_a^b \rho \frac{\partial u}{\partial t}(\rho, t) d\rho + \Phi(x, t), \quad 0 < a < x < b, 0 < t \leq T \tag{1.9}$$

together with (1.8). Let M and K be two positive integers and $h = \frac{b-a}{M}$, $\tau = \frac{T}{K}$. Denote

$$\begin{aligned} \Omega_h &= \{x_i \mid x_i = a + ih, 0 \leq i \leq M\}, & \Omega_\tau &= \{t_k \mid t_k = k\tau, 0 \leq k \leq K\}, \\ x_{i-1/2} &= (x_i + x_{i-1})/2, & t_{k-1/2} &= (t_k + t_{k-1})/2. \end{aligned}$$

If $u = \{u_i \mid 0 \leq i \leq M\}$ and $v = \{v_i \mid 0 \leq i \leq M\}$ are two mesh functions on Ω_h , take the notations

$$\begin{aligned} u_{i-1/2} &= (u_i + u_{i-1})/2, & \delta_x u_{i-1/2} &= (u_i - u_{i-1})/h, \\ \delta_x(x_i \delta_x u_i) &= (x_{i+1/2} \delta_x u_{i+1/2} - x_{i-1/2} \delta_x u_{i-1/2})/h \\ (u, v) &= h \sum_{i=1}^M x_{i-1/2} u_{i-1/2} v_{i-1/2}, & \|u\| &= \sqrt{(u, u)}, & \|u\|_\infty &= \max_{0 \leq i \leq M} |u_i|, \\ \|\delta_x u\| &= \sqrt{h \sum_{i=1}^M x_{i-1/2} (\delta_x u_{i-1/2})^2}, & \|\delta_x u\|_* &= \sqrt{h \sum_{i=1}^M \frac{1}{x_{i-1/2}} (\delta_x u_{i-1/2})^2}. \end{aligned}$$